

# Condensed Thoughts

Based on

\* ideas learned from collaboration colleagues

(Corclova, Dumitrescu, Thorngren)

\* more mathematical reframing by Gaiotto & Johnson-Freyd

\* detailed discussions (project?) with Freed & Hopkins

\* older ideas on finite homotopy theories w/ Freed, Hopkins, Lurie

↳ partially made precise by Scheimbauer-Walde

\* established ideas in (3D) physics / cond matter literature (Kitaev)

Theme: Calculus of topological defects in QFT

Emphasis: those coming from FHTs. (finite htpy theories)

Specific Goal: "Explain" Condensation of defects

or, what do we need in the target category to execute that?

"Deliverables": we isolated two ingredients for  $\dashv$

\* Dirichlet boundary condition (or, end of a defect)

\* Ostrik's principle: correspondence between

$\left. \begin{array}{l} \text{modules over algebras} \\ \text{in higher categories} \end{array} \right\} \longleftrightarrow \left\{ \text{Internal algebra objects} \right\}$

Applications [Recovering a pointed space from its symmetry TFT]

Recovering a  $(\pi$ -finite) space from its quantization ONLY POSSIBLE UP TO EM DUALITY

↳ (what is a quantum homotopy En type?)

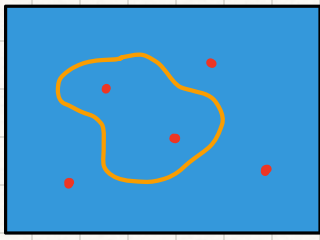
Moral: Filtration of defects by condensationsness

$\longleftrightarrow$  Postnikov filtration of the space

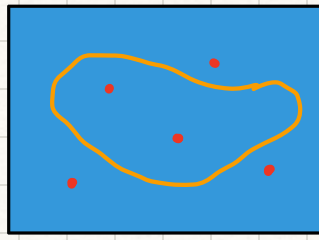
(CONDENSATION  
 $\Downarrow$   
CONNECTIVITY)

# Recap 1. Topological defects (extended operators)

\* can move freely in the vacuum



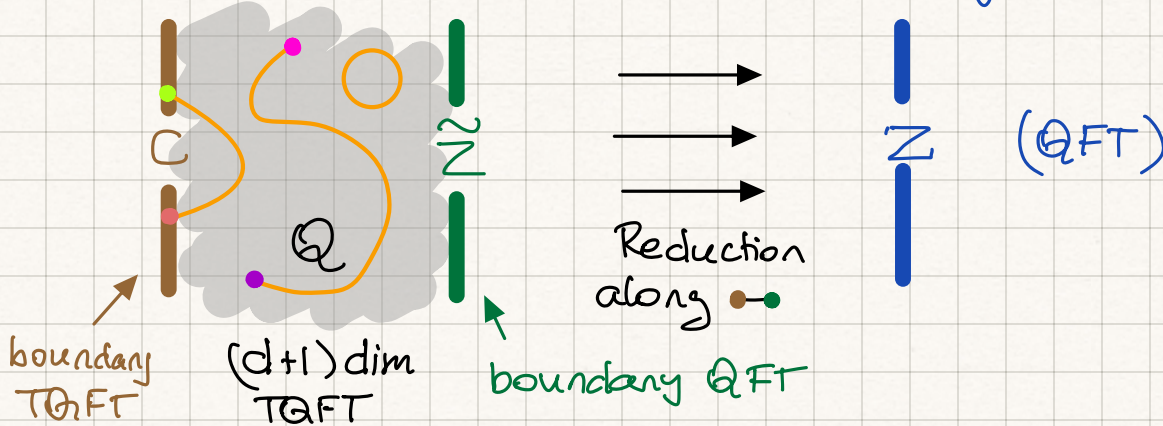
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topological operator

QFT operator

\* Our operators will be implemented by a quiche  $(Q, C)$



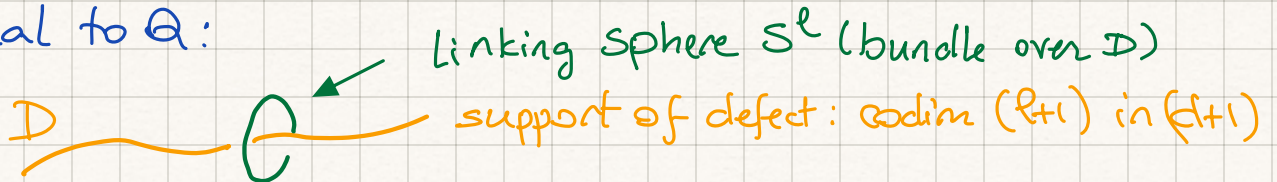
$(Q, C)$   
(crust)

Interpretation: Action of  $\text{End}_Q(C)$  on  $Z$ .

Mostly I'll be concerned with defects in the crust  $C$ .

# Recap 2. Calculus of defects

\* Internal to  $Q$ :



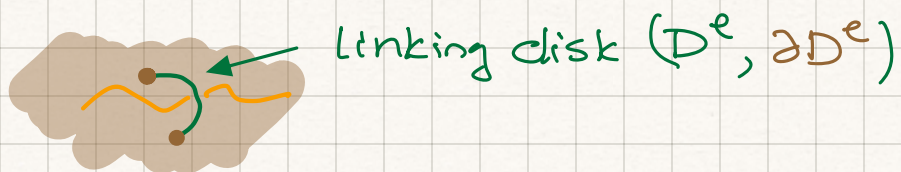
$Q(S^l)$ : defines a TQFT of dim  $(d+1-l)$

(Quantum) defect local label: boundary theory for  $Q(S^l)$   
= object in  $\text{Hom}_{\text{Set}}(\mathbb{1}^l, Q(S^l))$

Explanation: Remove a thickened  $D$ , place a locally (on  $D$ ) built state on the spherical boundary

Such states are built from boundary theories.

\* On the crust



$Q(D^e, \partial D^e)$ : TQFT of dim.  $(d+1-e)$ ... same story on labels

Example  $\mathcal{T}$ :  $(d+1)$ -category of  $d$ -algebras  $\text{Alg}(\text{Alg}(\text{Alg}(\dots)))$

If  $\mathcal{Q}$  is built from a  $\pi$ -finite space  $X$

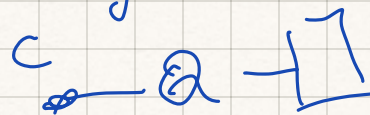
$C$  is the **canonical Dirichlet condition**

= one base-point in each component

$$Q(S^e) \stackrel{''}{=} \mathcal{Q}_{d-e}(\text{Map}(S^e; X))$$

= a  $(d-e)$ -algebra built by iterated crossed products with the  $\pi_i \text{Map}(S^e; X)$  in descending order

$$Q(D^e, \partial D^e) \stackrel{''}{=} \bigoplus_{\pi_0 X} \mathcal{Q}_{d-e}(\Omega^e X)$$



Boundary theories for these are modules in  $(d-e-1)$  algebras

$\equiv (d-e)$  dim. TQFTs with actions of  $\Omega \text{Map}(S^e; X), \dots$

Some can be obtained from spaces  $Y \rightarrow \text{Map}(S^e; X)$  by quantizing the **homotopy fiber**

(The loop space of the base acts on the homotopy fiber)

Addenda:  $X$  may carry a generalized Dijkgraaf-Witten twist.

Caution (orientations)

The spheres rotate along the defect.

the  $SO(\ell+1)$  action on  $\text{Map}(S^\ell; X)$  may be nontrivial

$\Rightarrow$  care needed for tangential structures

# Physicist's wish list for a defect • of codimension $\ell$

• • • • •  
many defects

↓  
zoom out  
and squint  
really hard



defect of codimension  
 $(\ell-1)$ .

More seriously:

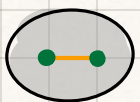
- \* Condensed defects should be **endable**,
- \* A closed condensed defect can be **ripped open** and **closed back up**
- \* This allows any **linking defect** to slide away.
- \* Iterating: The support of  $k$ -fold condensed defects may be shrunk by  $k$  dimensions.

This phenomenon accounts for claims of a "calculus of defects" that mixes dimensions.

(Some of the defects in the product may be shrunk).

Interpretation [Loose with framings and duals!]

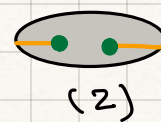
Endable defects of codim  $\ell \rightarrow$  algebra objects in defects codim  $(\ell+1)$



algebra in one  
lower dimension



But,  $m$  need not be an isomorphism so need a condition on the **end** to recover



This is the **(strong) Dirichlet condition**

Example At the vector space level (**line**, **end**) we need a collection of ends so that the (2) contain (1) in their span.

## Dirichlet condition on a boundary theory

$(d+1)$ -dim. TQFT  $\mathcal{Q}$  (e.g.  $\mathcal{Q}_d(x)$ )  $\leftrightarrow \mathcal{Q}(pt) \in \mathcal{T}, (d+1)$  cat.

Boundary theory  $\mathcal{C} \leftrightarrow \beta := \mathcal{C}(pt) \in \text{Hom}_{\mathcal{T}}(\mathbb{1}, \mathcal{Q}(pt))$

$\mathcal{C}$  is Dirichlet iff:

(i) The  $\text{End}_{\mathcal{T}}(\mathbb{1})$ -module categories  $\text{Hom}_{\mathcal{T}}(\mathbb{1}, \mathcal{Q}(pt))$  and  $\text{Mod-}\beta^R \circ \beta$  are equivalent (Oshikawa correspondence)

(" $\beta$  generates all boundary theories for  $\mathcal{Q}$ ")

(ii)  $\mathcal{Q}(pt)$  is local at  $\mathbb{1}$ : For  $\forall \gamma \in \text{Ob } \mathcal{T}$ , have an equivalence

$\text{Hom}(\mathcal{Q}(pt), \gamma) \rightarrow \text{Hom}_{\text{End}(\mathbb{1})}(\text{Hom}(\mathbb{1}, \mathcal{Q}(pt)), \text{Hom}(\mathbb{1}, \gamma))$

(" $\mathcal{Q}$  is equivalent to the TQFT generated by  $\beta^R \circ \beta$ ")

Example:  $\mathcal{T} = 2$ -category of (Alg, Bimod, Intertwiners)

$\gamma$ - $\mathcal{Q}$  bimodules  $\leftrightarrow$  (co)continuous functors  $\mathcal{Q}\text{-Mod} \rightarrow \gamma\text{-Mod}$ .

Remark: It should be the case that (i) and (ii) can be reworded in the language of retracts as in Gaiotto & Johnson-Freyd

as  $\text{Id}_{\mathcal{Q}}$  is a retract of  $\beta \circ \beta^t \in \text{End}(\mathcal{Q})$ .

## Dirichlet condition on the end of a defect

Treat the defect as an embedded TQFT and repeat.



condenser, codim  $(l+1)$   
algebra defect



condensate, codim  $l$

Quasi-theorem\* A defect with a Dirichlet end may be cut open and closed up again.

\*provable after making more precise

## Examples of Dirichlet conditions

\* FHT:  $Y \rightarrow X$  is a Dirichlet condition for dimension  $(d+1)$

$$\Leftrightarrow \begin{cases} \pi_0 Y \rightarrow \pi_0 X \text{ onto} \\ \pi_d Y \rightarrow \pi_d X \text{ zero}^* \end{cases} \quad \begin{array}{l} \updownarrow \text{EM dual.} \\ \text{(*exceptions if } d=2 \text{ \& } \pi_1 X \text{ nonab.)} \end{array}$$

eg:  $d=1$ ,  $X = BG$   $Y = BH$

Dirichlet  $\mathcal{D}$ : a representation containing every irrep of  $G$ .

If  $G$  abelian, only happens for  $\text{Ind}_H^G$  if  $H \rightarrow G$  is trivial.

\* In particular, if  $X$  is connected,  $(pt) \rightarrow X$  is Dirichlet.

(**canonical** Dirichlet condition:

the homotopy fiber is  $\Omega X$  with its translation action).

\* Dirichlet ends for FHT defects:

$Y \rightarrow \Omega^e X$  must satisfy the  $\pi_0, \pi_{d-e}$  as above

\* Fusion categories  $\mathcal{F}$

$\mathcal{F}$  indecomposable  $\Rightarrow$  any nonzero module cat  $\mathcal{M}$  is Dirichlet

$$\mathcal{F} \stackrel{\mathcal{M}}{\cong} \text{End}_{\mathcal{F}}(\mathcal{M}).$$

$$\mathcal{F} \cong \mathcal{F}_1 \oplus \mathcal{F}_2, \quad \mathcal{M} = \mathcal{F}_1 \oplus \mathcal{O} \quad - \text{not Dirichlet.}$$

In particular, for gauge theory in  $\dim \geq 3$

the Neumann boundary theory is Dirichlet.

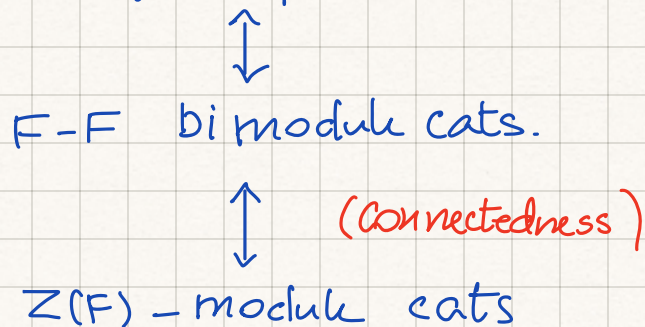
# Example: Condensation in 3D theories

Theorem In a simple Turaev-Viro or Reshetikhin-Turaev theory:

- (1) only the transparent line operator is condensed
- (2) all surface operators are condensed.

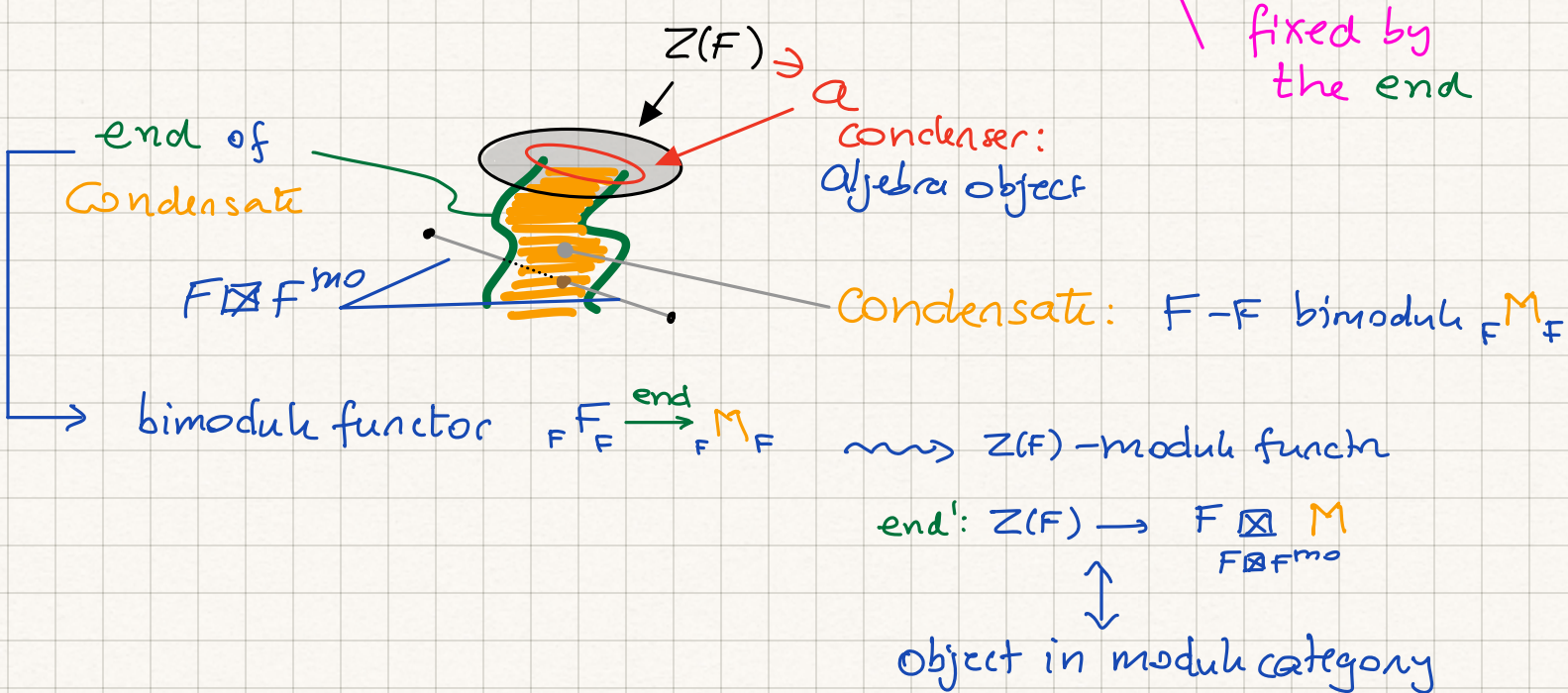
PF (1) Only scalar point operators (simple; connected)

(2) Surface operators



algebra objects in  $Z(F)$  up to Morita equivalence

fixed by the end



Remark In RT theories, interfaces are " $Z(F)$ "-modules by folding.

\* Key (Non)Fact

The Dirichlet property is **NOT** preserved by dimensional reduction.

Eg: \* 3D gauge theory reduced on  $S^1$ :

No  $\mathbb{Z}$  condition generates the Drinfeld center upon reduction.

\*  $\text{Maps}(M; *) \rightarrow \text{Map}(M; X)$  is **NOT** Dirichlet

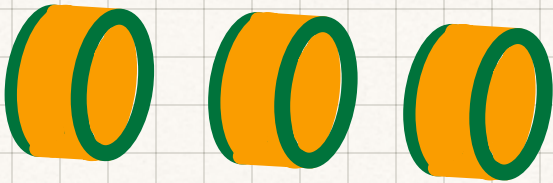
← usually not connected!

Key problem: We can condense from  $\text{dim. } (d-e)$  to  $(d-e+1)$

But then we seem stuck because our 'end' is no longer Dirichlet:



But

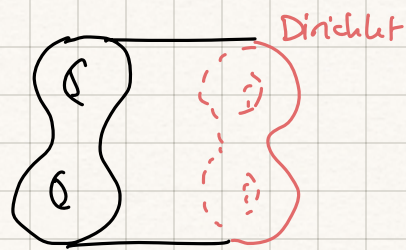


seems not condensable?

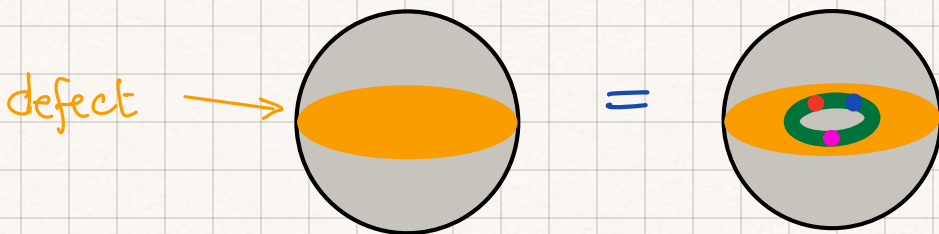
Theorem (FHT case but probably in general)

The dimensional reduction of a Dirichlet end can be made Dirichlet after embedding additional self defects.

Example We can span the space of states for a  $d$ -manifold by inserting higher codim defects in the Dirichlet boundary.



Theorem At the top dimension (if  $d+1 \geq 3$ ) we can 'rip open' the defect without changing correlators by embedding suitable self-defects in the Dirichlet end.



same vectors in the space for the (defective) sphere ( $S^2$ , equator)



Classical Condensation: Crust defects in the  $(d+1)$  dim Quiche from a space  $X$  with its Dirichlet (basepoint) condition

$\mathcal{Q}_m(Y) :=$  quantization as algebra in an  $m$ -category

Codimension of defects	Local quantum labels are modules over...	Contributing homotopy groups, placed in $\dim [k]$	
		all/ <span style="border: 1px solid red;">new</span> defects	Condensed
$d$ (pts)	$\mathcal{Q}_0(\Omega^d X)$	<span style="border: 1px solid red;"><math>\pi_d[0]</math></span>	None
$d-1$ (lines)	$\mathcal{Q}_1(\Omega^{d+1} X)$	<span style="border: 1px solid red;"><math>\pi_{d-1}[0]</math></span> , $\pi_d[1]$	$\pi_d[1]$
$d-2$ (surfaces)	$\mathcal{Q}_2(\Omega^{d-2} X)$	<span style="border: 1px solid red;"><math>\pi_{d-2}[0]</math></span> , $\pi_{d-1}[1]$ , $\pi_d[2]$	$\pi_{d-1}[1]$ , $\pi_d[2]$
$k$	$\mathcal{Q}_{d-k}(\Omega^k X)$	<span style="border: 1px solid red;"><math>\pi_k[0]</math></span> , $\pi_{k+1}[1]$ , $\pi_{k+2}[2]$ , ...	$\pi_{k+1}[1]$ , $\pi_{k+2}[2]$ , ...
$2$	$\mathcal{Q}_{d-2}(\Omega X)$	<span style="border: 1px solid red;"><math>\pi_2[0]</math></span> , $\pi_3[1]$ ..., $\pi_d[d-2]$	$\pi_3[1]$ ..., $\pi_d[d-2]$
$1$ (interface)	$\mathcal{Q}_{d-1}(X)$	<span style="border: 1px solid red;"><math>\pi_1[0]</math></span> , $\pi_2[1]$ , ..., $\pi_d[d-1]$	$\pi_2[1]$ , $\pi_3[2]$ , ..., $\pi_d[d-1]$

"Classically" condensed defects are those in positive htpy degree. High Postnikov fibers are increasingly condensed.

Classical shrinking of support: maps from manifolds to highly connected spaces can be supported in high co-dimension.

The listing also suggests that  $[\mathcal{Q}_d(X), \mathcal{Q}_{d-1}(\Omega X)]$  determines  $(X, *)$ .

That is the case, and actions of this quiche on a QFT  $\Leftrightarrow$  actions of  $\Omega X$ .

## Shrinking of support

Classical:

Assume  $Z$  connected,  $\pi_1 Z$  abelian, acting trivially on  $\pi_n Z$ .

Then:  $\pi_k$  Based Map  $(M; Z) \cong \pi_k \text{Map}(M, Z)$  whenever  $\pi_k Z = 0$ .

$M$  manifold  $\Rightarrow$  Based Map can be retracted to codim 1.

$Z$  highly connected  $\Rightarrow$  Can shrink the support of maps without damaging the low homotopy groups.

Quantum:

Even if  $X$  is highly connected, high-dimensional defects need not be condensed: because you could stack a TFT w/o boundary conditions on top of the defect. But that's the only problem:

Theorem If the quantum defect label admits a Dirichlet  $\partial$ , and maps to the  $> 0$  part of  $\Omega^k X$ , then the defect is condensed and its support may be retracted to codim 1.

may be theorem one day

~~Alt Theorem If the  $\Omega^k X$  action on the label is induced from a high Postnikov coin, then in top dimension the defect is a constant TFT with an embedded defect.~~

Repeating the bookkeeping for  $\mathcal{G}_d(X)$  alone "doubles" the homotopy groups, because  $\text{Maps}(S^d; X)$  replace  $\Omega^d X$ .

For codim  $(k+1)$  defects, space is  $\text{Map}(S^k; X)$ ;

$$\begin{cases} \pi_0, \dots, \pi_{d-k} & \text{from } X \\ \pi_k[0], \dots, \pi_d[d-k] & \text{from } \Omega^k X \end{cases}$$

Positions  $p$  and  $(d-p)$  of homotopy groups are indistinguishable

Appearance of  $\pi_k \text{Sol}$  for  $\mathcal{Q}_{d-k}(\text{Map}(S^k; X))$  can't be distinguished from occurrence of  $\pi_k^\vee$  in deg.  $(d-k)$ .

so  $\pi_k$  will occur in degrees  $k$  and  $(d-k)$  on the list

Expected

Theorem

The quantization  $\mathcal{Q}_d(X)$  determines  $X$  "up to electro-magnetic duality".

Morally:  $(d+1)$ -dim quantization "folds"

homotopy  $d$ -types (with Dijkgraaf-Witten twist)

$$\text{via } \pi_k \longleftrightarrow \pi_{d-k}$$

? Answer Homotopy types with "poly-k-invariants".

$X$  has:  $\pi_i(X)$ ,  $k \in \bigoplus_p C^p(X; \pi_{p-1}X)$  ( $+ \tau \in H^{d+1}(X; \mathbb{C}^*)$ )

Morally: cochains on  $X =$  functions on the structured set  $\bigoplus \pi_p(X)$

$\Rightarrow (k, \tau)$  is a function on  $\bigoplus \pi_p(X) \times \bigoplus \pi_p(X)^\vee$   
"X" "X $^\vee$ "

- Linear on  $X^\vee$
- Of total degree  $(d+1)$ , if you shift  $X^\vee$  to degree  $(-d)$

EM duality switches  $X$  and  $X^\vee$ .

(That's why we can define it for linear  $k \iff$   $\infty$  loop space  $X$ ).

To extend this we need nonlinear (in  $X^\vee$ )  $k$ -invariants

Proposition This makes sense in rational homotopy and defines  $E_d$  structures on  $C^*(X; \mathbb{Q})$ .

No comparable statement for  $\pi$ -finite spaces

General procedure TQFT  $\mathcal{Q}$ , target cat.  $\mathcal{J}$

Have an adjunction

$$\text{Hom}_{\Omega^{l+1} \mathcal{J}}(\mathbb{1}^{l+1}, \mathcal{Q}(S^{l+1})) \cong \text{End}_{\text{Hom}(\mathbb{1}^l, \mathcal{Q}(S^l))}(\mathcal{Q}(D^{l+1}))$$

$$\begin{array}{ccc} \mathbb{1}^l & \xrightarrow{\mathcal{Q}(D^{l+1})} & \mathcal{Q}(S^l) \\ & \nwarrow \mathcal{Q}(D^{l+1})^* & \uparrow \\ & & \mathcal{Q}(D^{l+1})^* \circ \mathcal{Q}(D^{l+1}) \\ & & \mathcal{Q}(S^l) \end{array}$$

$$\Omega_{\mathcal{J}}^l := \underbrace{\text{End}(\text{End}(\dots(\text{End}(\mathbb{1}_0)\dots))}_{l \text{ times}}$$

Ostrik principle:

algebra objects in LHS = RHS  $\longleftrightarrow$  modules over RHS (with a generator)

Localization Condition (from def of Dirichlet end)

Dirichlet-endable objects in  $\text{Hom}(\mathbb{1}^l, \mathcal{Q}(S^l))$  are determined by their "localization at the unit"  $\mathcal{Q}(D^{l+1})$   
 Localization at  $\mathcal{Q}(D^{l+1})$  is an equivalence

$$\begin{array}{l} \mathcal{M} \in \text{Hom}_{\Omega_{\mathcal{J}}^l}(\mathbb{1}^l, \mathcal{Q}(S^l)) \\ + \text{ (Dirichlet) } \\ \text{end} \in \text{Hom}(\mathcal{Q}(D^{l+1}), \mathcal{M}) \end{array}$$

$$\xrightarrow{\text{localize}} \begin{array}{l} \text{Hom}_{\Omega_{\mathcal{J}}^l}(\mathcal{Q}(D^{l+1}), \mathcal{M}) \\ \text{End}_{\text{Hom}(\mathbb{1}^l, \mathcal{J}(S^l))}(\mathcal{Q}(D^{l+1})) \end{array}$$

algebra object

$$\mathcal{A} \in \text{Hom}(\mathbb{1}^{l+1}, \mathcal{Q}(S^{l+1}))$$

$$\cong \begin{array}{l} + \text{ generator} \\ \longleftarrow \\ \text{Ostrik} \end{array}$$

Dirichlet condition on end says that you can "condense back" to the original (defect, end).

## Overall

$\mathcal{A}$  = algebra in  $\text{codim}(\ell+2)$  local defects  $\text{Hom}(\mathbb{1}_{\ell+1}, \mathcal{G}_d(x)[s^{\ell+1}])$   
+ regular module

→  $\text{End} \dots [\mathcal{A}(D^{\ell+1})]$  module (+ generator)

→  $\text{codim}(\ell+1)$  defect label in  $\text{Hom}(\mathbb{1}^{\ell}, \mathcal{G}_d(x)[s^{\ell}])$   
(+ defect end) (condensate)